


## Center of Mass



Diver's motion is pure translation


it is translation plus rotation


There is one point that moves in same path a particle would take if subjected to same force as diver This point is called center of mass (CM)

## Center of Mass

For two particles center of mass reads

$$
x_{C M}=\frac{m_{A} x_{A}+m_{B} x_{B}}{m_{A}+m_{B}}=\frac{m_{A} x_{A}+m_{B} x_{B}}{M}
$$

Where $M$ is Total Mass


## Center of Mass

We can generalize from two particles in one dimension to a system of $N$ particles in three dimensions

For a continuous distribution of mass

$$
\vec{r}_{c m}=\frac{1}{M} \int \vec{r} d m
$$

$$
\begin{aligned}
& \text { ensions dimension } \quad \begin{aligned}
& M x_{c m}=\sum_{i}^{N} m_{i} x_{i} \\
& M y_{c m}=\sum_{i}^{N} m_{i} y_{i} \\
& M z_{c m}=\sum_{i}^{N} m_{i} z_{i} \\
& M=\sum_{i}^{N} m_{i} \\
& M \vec{r}_{c m}=\sum_{i}^{N} m_{i} \vec{r}_{i} \\
& \underbrace{x}_{c m}=x_{c m} \hat{\imath}+y_{c m} \hat{\jmath}+z_{c m} \hat{k}
\end{aligned}
\end{aligned}
$$

## Center of Mass of a uniform semicircular hoop

Use polar coordinates

$$
y=r \sin \theta
$$

Distance of points on semicircle from origin is $r=R$

$\vec{r}_{c m}=\frac{1}{M} \int(x \hat{\imath}+y \hat{\jmath}) d m=\frac{1}{M} \int R(\cos \theta \hat{\imath}+\sin \theta \hat{\jmath}) d m$
mass per unit length is $\lambda=d m / d s$

$$
\begin{gathered}
d m=\lambda d s=\lambda R d \theta \\
\vec{r}_{c m}=\frac{1}{M} \int_{0}^{\pi} R(\cos \theta \hat{\imath}+\sin \theta \hat{\jmath}) \lambda R d \theta
\end{gathered}
$$

## Center of Mass of a uniform semicircular hoop (cont'd)

$$
\begin{aligned}
& \vec{r}_{c m}=\frac{\lambda R^{2}}{M}\left(\hat{\imath} \int_{0}^{\pi} \cos \theta d \theta+\hat{\jmath} \int_{0}^{\pi} \sin \theta d \theta\right) \\
& \begin{array}{c}
x=R \cos \theta \\
d m=\lambda d s=\lambda R d \theta \\
d i n \\
\vec{r}_{c m}
\end{array} \\
&=\frac{R}{\pi}\left(\hat{\imath} \int_{0}^{\pi} \cos \theta d \theta+\hat{\jmath} \int_{0}^{\pi} \sin \theta d \theta\right) \\
&=\frac{R}{\pi}\left(\left.\hat{\imath} \sin \theta\right|_{0} ^{\pi}-\left.\hat{\jmath} \cos \theta\right|_{0} ^{\pi}\right) \\
&=\frac{2}{\pi} R \hat{\jmath}
\end{aligned}
$$

Curiously it is outside material of semicircular hoop!

## Center of Mass for human body

High jumpers have developed a technique where their CM actually passes under bar as they go over it


This allows them to clear higher bars

## Center of Mass \& Translational Motion

- Sum of all forces acting on a system is equal to total mass of system multiplied by acceleration of center-of-mass
- For each internal force acting on a particle in system there is an equal and opposite internal force acting on some other particle of system

$$
M \vec{a}_{C M}=\vec{F}_{n e t, e x t}
$$

## Changing Places in a Rowboat

Pete (mass 80 kg ) and Dave (mass 120 kg ) are in a rowboat (mass 60 kg ) on a calm lake Dave is near bow of boat, rowing, and Pete is at stern, 2 m from Dave Dave gets tired and stops rowing
Pete offers to row, so after boat comes to rest they change places
How far does boat moves as Pete and Dave change places

$M X_{\mathrm{CM}_{i}}=m_{\text {Pete }} x_{\text {Pete }_{i}}+m_{\text {Dave }} x_{\mathrm{Dave}_{i}}+m_{\text {boat }} x_{\mathrm{boat}_{i}}$
$M X_{\mathrm{CM}_{f}}=m_{\text {Pete }} x_{\text {Pete }_{f}}+m_{\text {Dave }} x_{\text {Dave }_{f}}+m_{\text {boat }} x_{\text {boat }_{f}}$ $M \Delta X_{\mathrm{CM}}=m_{\text {Pete }} \Delta x_{\text {Pete }}+m_{\text {Dave }} \Delta x_{\text {Dave }}+m_{\text {boat }} \Delta x_{\text {boat }}$

$$
M a_{\mathrm{CM}}=F_{\mathrm{ext}, \mathrm{net}}=\sum_{i} F_{\mathrm{ext}_{i}}=0
$$

$0=m_{\text {Pete }}\left(\Delta x_{\text {boat }}+L\right)+m_{\text {Dave }}\left(\Delta x_{\text {boat }}-L\right)+m_{\text {boat }} \Delta x_{\text {boat }}$
$\Delta x_{\text {boat }}=\frac{L\left(m_{\text {Dave }}-m_{\text {Pete }}\right)}{m_{\text {Dave }}+m_{\text {Pete }}+m_{\text {boat }}}=0.31 \mathrm{~m}$

$x_{\text {Davef }}$

## Center of Mass Work

For a system of particles we have

$$
\begin{aligned}
\vec{F}_{\mathrm{net}, \mathrm{ext}} & =\sum \vec{F}_{\mathrm{iext}}=M \vec{a}_{\mathrm{cm}} \\
M & =\sum_{i} m_{i} \\
\vec{F}_{\mathrm{net}, \mathrm{ext}} \cdot \vec{v}_{\mathrm{cm}} & =M \vec{a}_{\mathrm{cm}} \cdot \vec{v}_{\mathrm{cm}}=\frac{d}{d t}\left(\frac{1}{2} M v_{\mathrm{cm}}^{2}\right)=\frac{d K_{\mathrm{trans}}}{d t}
\end{aligned}
$$

Center of mass work - translational kinetic energy relation

$$
\int_{1}^{2} \vec{F}_{\mathrm{net}, \mathrm{ext}} \cdot d \vec{l}_{c m}=\Delta K_{\mathrm{trans}}
$$



In SI system, units of power are watts:

$$
1 W=1 \mathrm{~J} / \mathrm{s}
$$

Difference between walking and running up these stairs is power

- change in gravitational potential energy is same-


## Conservation Theorems: Momentum



Luis Anchordoqui

## Momentum of a particle

Originally introduced by Newton as quantity of motion

$$
\text { momentum is defined as } \vec{p}=m \vec{v}
$$

Using Newton"s second law we can relate momentum of a particle to force acting on particle

$$
\frac{d \vec{p}}{d t}=\frac{d(m \vec{v})}{d t}=m \frac{d \vec{v}}{d t}=m \vec{a}
$$

Substituting force $\vec{F}_{\text {net }}$ by

$$
\vec{F}_{\mathrm{net}}=\frac{d \vec{p}}{d t}
$$

Net force acting on a particle equals time rate of change of particle"s momentum
In his famous treatise Principia (1687) Newton presents second law of motion in this form

## Conservation of Momentum

Total momentum of a system of particles reads

$$
\begin{gathered}
\vec{P}_{\mathrm{sys}}=\sum_{i} m_{i} \vec{v}_{i}=\sum_{i} \vec{p}_{i} \\
\vec{P}_{\mathrm{sys}}=\sum_{i} m_{i} \vec{v}_{i}=M \vec{v}_{\mathrm{cm}} \\
M=\sum_{i} m_{i} \\
\frac{d \vec{P}_{\mathrm{sys}}}{d t}=M \frac{d \vec{v}_{\mathrm{cm}}}{d t}=M a_{\mathrm{cm}}
\end{gathered}
$$

According to Newton" s second law

$$
\begin{aligned}
& \quad \sum_{i} \vec{F}_{\mathrm{ext}}=\vec{F}_{\mathrm{net}, \mathrm{ext}}=\frac{d \vec{P}_{\mathrm{sys}}}{d t} \\
& \text { If } \sum_{\text {If }} \vec{F}_{\mathrm{ext}}=0 \text { then } \vec{P}_{\mathrm{sys}}=M \vec{v}_{\mathrm{cm}}=\mathrm{constant} \\
& \text { If sum of external forces on a system remains zero } \\
& \text { total momentum of system is conserved }
\end{aligned}
$$

## A runaway railroad car

A runaway $14,000 \mathrm{~kg}$ railroad car is rolling horizontally at $4 \mathrm{~m} / \mathrm{s}$ toward a switchyard
As it passes by a grain elevator 2000 kg of grain suddenly drops into the car
How long does it take car to cover 500 m distance from elevator to switchyard?

Assume that grains falls straight down and that slowing due to rolling friction or air drag is negligible


## A runaway railroad car

$$
\begin{aligned}
& \sum_{i} \vec{F}_{\mathrm{ext}_{i}}=\vec{F}_{g_{\mathrm{grain}}}+\vec{F}_{g_{\mathrm{car}}}+\vec{F}_{n}=\frac{d \vec{P}_{\mathrm{sys}}}{d t} \\
& \quad \times \text { component of net external force is zero }
\end{aligned}
$$



$$
\begin{gathered}
P_{\mathrm{sys}, \mathrm{x}_{i}}=P_{\mathrm{sys}, \mathrm{x}_{f}} \\
\left(m_{\mathrm{car}}+m_{\mathrm{grain}}\right) v_{x_{f}}=m_{\mathrm{car}} v_{x_{i}}
\end{gathered}
$$



Before

$$
v_{x_{f}}=\frac{m_{\mathrm{car}} v_{x_{i}}}{m_{\mathrm{car}}+m_{\mathrm{grain}}} \Rightarrow \Delta t=\frac{d}{v_{x_{f}}}=143 \mathrm{~s}
$$

## Collisions and impulse

When two objects collide
they usually exert very large forces on each other for very brief time Impulse of a force exerted during a time $\Delta t=t_{f}-t_{i}$ is a vector defined as


$$
\vec{I}=\int_{t_{i}}^{t_{f}} \vec{F} d t
$$

Impulse is a measure of both
strength and duration of collision force

$$
\vec{I}=\int_{t_{i}}^{t_{f}} \vec{F} d t=\int_{t_{i}}^{t_{f}} \frac{d \vec{p}}{d t} d t=\vec{p}_{f}-\vec{p}_{i}
$$

Impulse momentum theorem for a particle

$$
\vec{I}_{\mathrm{net}}=\Delta \vec{P}
$$

Impulse momentum theorem for a system

$$
\vec{I}_{\mathrm{net}, \mathrm{ext}}=\int_{t_{i}}^{t_{f}} \vec{F}_{\mathrm{net}, \mathrm{ext}} d t=\Delta \vec{P}_{\mathrm{sys}}
$$

## Average force

Since time of collision is very short we need not worry about exact time dependence of force and can use average force


## Conservation of Energy and Momentum in Collisions

## Momentum is conserved in all collisions

Collisions in which kinetic energy is conserved as well are called elastic collisions and those in which it is not are called inelastic


If elastic


Collision



If inelastic

## Perfectly Elastic and Perfectly Inelastic Collisions

Conservation of momentum

$$
m_{1} v_{1 \mathrm{f}}+m_{2} v_{2 \mathrm{f}}=m_{1} v_{1 \mathrm{i}}+m_{2} v_{2 \mathrm{i}}
$$

In elastic collisions kinetic energy of system is conserved

$$
\begin{aligned}
\frac{1}{2} m_{1} v_{1 \mathrm{f}}^{2}+\frac{1}{2} m_{2} v_{2 \mathrm{f}}^{2} & =\frac{1}{2} m_{1} v_{1 \mathrm{i}}^{2}+\frac{1}{2} m_{2} v_{2 \mathrm{i}}^{2} \\
m_{2}\left(v_{2 \mathrm{f}}^{2}-v_{2 \mathrm{i}}^{2}\right) & =m_{1}\left(v_{1 \mathrm{i}}^{2}-v_{1 \mathrm{f}}^{2}\right) \\
m_{2}\left(v_{2 \mathrm{f}}-v_{2 \mathrm{i}}\right)\left(v_{2 \mathrm{f}}+v_{2 \mathrm{i}}\right) & =m_{1}\left(v_{1 \mathrm{i}}-v_{1 \mathrm{f}}\right)\left(v_{1 \mathrm{i}}+v_{1 \mathrm{f}}\right)
\end{aligned}
$$

From conservation of momentum

$$
m_{2}\left(v_{2 \mathrm{f}}-v_{2 \mathrm{i}}\right)=m_{1}\left(v_{1 \mathrm{i}}-v_{1 \mathrm{f}}\right)
$$

Taking ratio of these two equations

$$
\left(v_{2 \mathrm{f}}+v_{2 \mathrm{i}}\right)=\left(v_{1 \mathrm{f}}+v_{1 \mathrm{i}}\right)
$$

Rearranging we obtain relative velocities in an elastic collision

$$
\left(v_{1 \mathrm{i}}-v_{2 \mathrm{i}}\right)=\left(v_{2 \mathrm{f}}-v_{1 \mathrm{f}}\right)
$$

In perfectly inelastic collisions objects have same velocity after collision (often because they stuck together)

$$
v_{1 \mathrm{f}}=v_{2 \mathrm{f}}=v_{\mathrm{cm}}
$$

A meteor whose mass was about $10^{8} \mathrm{~kg}$ struck Earth with a speed of about $15 \mathrm{~km} / \mathrm{s}$ and came to rest in Earth
(a) What was Earth's recoil speed?
(b) What fraction of meteor's kinetic energy was transformed to kinetic energy of Earth?
(c) By how much did Earth kinetic energy change as a result of this collision?


$$
M_{\oplus}=6 \times 10^{24} \mathrm{~kg}
$$

(a) $\quad m_{\text {meteor }} v_{\text {meteor }}=\left(m_{\text {meteor }}+M_{\oplus}\right) v_{f}$


$$
v_{f}=6 \times 10^{-11} \mathrm{~m} / \mathrm{s}
$$

(b)

$$
\frac{K_{f}^{\text {Earth }}}{K_{i}^{\text {meteor }}}=8 \times 10^{-15}
$$

(c)

$$
\Delta K_{\mathrm{Earth}}=K_{f}^{\mathrm{Earth}}-K_{i}^{\mathrm{Earth}}=\frac{1}{2} M_{\oplus} v_{f}^{2}=10,800 \mathrm{~J}
$$

## Elastic Collisions in One Dimension

Here we have two objects colliding elastically
We know masses and initial speeds

$x$
Since both momentum and kinetic energy are conserved we can write two equations This allows us to solve for two unknown final speeds


## Elastic collision of a neutron and a nucleus

A neutron of mass $m_{\mathrm{n}}$ and speed $v_{\mathrm{ni}}$ undergoes a head-on elastic collision with a carbon nucleus of mass $m_{C}$ initially atrest
(a) What are final velocities of both particles?
(b)What fraction $f$ of its initial kinetic energy does neutron lose?


$$
\begin{aligned}
& v_{\mathrm{ni}} \\
& \text { for elastic collision } v_{C_{f}}-v_{n_{f}}=v_{n_{i}} \Rightarrow v_{C_{f}}=v_{n_{i}}+v_{n_{f}} v_{n_{i}}=m_{n} v_{n_{f}}+m_{C} v_{C_{f}} \\
& m_{n} v_{n_{i}}=m_{n} v_{n_{f}}+m_{C}\left(v_{n_{i}}+v_{n_{f}}\right) \\
& v_{n_{f}}=-\frac{m_{C}-m_{n}}{m_{n}+m_{C}} v_{n_{i}} \\
& v_{C_{f}}=v_{n_{i}}-\frac{m_{C}-m_{n}}{m_{n}+m_{C}} v_{n_{i}}=\frac{2 m_{n}}{m_{n}+m_{C}} v_{n_{i}} \\
& f=-\frac{\Delta K_{n}}{K_{n_{i}}}=\frac{K_{C_{f}}}{K_{n_{i}}}=\frac{m_{C}}{m_{n}}\left(\frac{v_{C_{f}}}{v_{n_{i}}}\right)^{2}=\frac{4 m_{n} m_{C}}{\left(m_{n}+m_{C}\right)^{2}}
\end{aligned}
$$

## Collisions in Two or Three Dimensions

Conservation of energy and momentum can also be used to analyze collisions in two or three dimensions but unless situation is very simple math quickly becomes unwieldy

If moving object collides with an object initially at rest
knowing masses and initial velocities is not enough we need to know angles as well in order to find final velocities


A novice pool player is faced with corner pocket shot shown in figure Relative dimensions are also given
Should player be worried about this being a "scratch shot" in which cue ball will also fall into a pocket?


Give details


In elastic collision between two objects of equal mass with target at rest angle between final velocities of objects is $90^{\circ}$

Momentum conservation

$$
m \vec{v}=m \vec{v}_{A}^{\prime}+m \vec{v}_{B}^{\prime} \Rightarrow \vec{v}=\vec{v}_{A}^{\prime}+\vec{v}_{B}^{\prime}
$$



Kinetic energy conservation $\frac{1}{2} m v^{2}=\frac{1}{2} m v_{A}^{\prime 2}+\frac{1}{2} m v_{B}^{\prime 2} \Rightarrow v^{2}=v_{A}^{\prime 2}+v_{B}^{\prime 2}$
Applying law of cosines $-v^{2}=\vec{v}_{A}^{\prime 2}+\vec{v}_{B}^{\prime 2}-2 \vec{v}_{A}^{\prime} \vec{v}_{B}^{\prime} \cos \theta$
Equating two expressions for $v^{2}$ leads to

$$
\begin{aligned}
& v_{A}^{\prime 2}+v_{B}^{\prime 2}=v_{A}^{\prime 2}+v_{B}^{\prime 2}-2 v_{A}^{\prime} v_{B}^{\prime} \cos \theta \\
& \cos \theta=0 \Rightarrow \theta=90^{\circ}
\end{aligned}
$$



Assume that target ball is hit correctly so that it goes in pocket
From geometry of right triangle

$$
\theta_{1}=\arctan [1 / \sqrt{3}]=30^{\circ}
$$

From geometry of left triangle

$$
\theta_{2}=\arctan [3 / \sqrt{3}]=60^{\circ}
$$

Because balls will separate at $90^{\circ}$. if target ball goes in pocket this does appear to be a good possibility of a scratch shot

