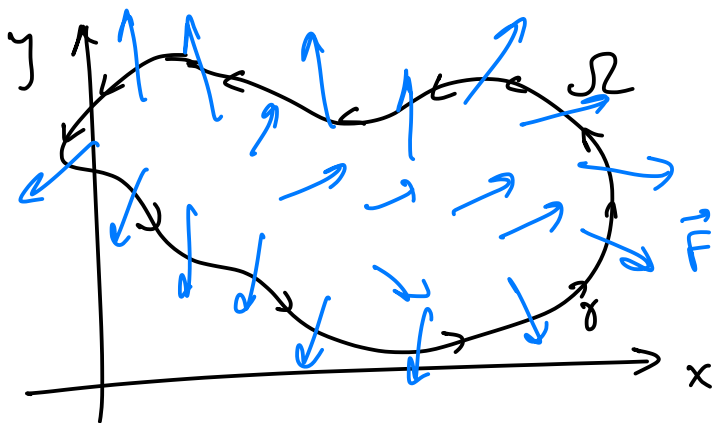


# Generalizations of Green's Theorem

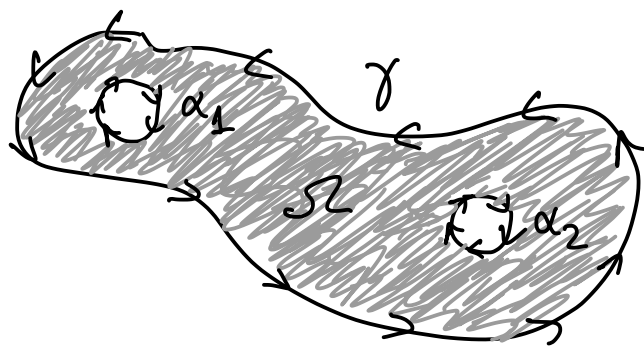
Recall Green's Theorem:

$\Omega \subset \mathbb{R}^2$  simply-connected  
 $\vec{F}: \Omega \subset \mathbb{R}^2 \rightarrow \mathbb{R}^2$  vector field  
 $\vec{F} = (M, N) = M\hat{i} + N\hat{j}$



$$\int_{\gamma} \vec{F} \cdot d\gamma = \iint_{\Omega} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$$

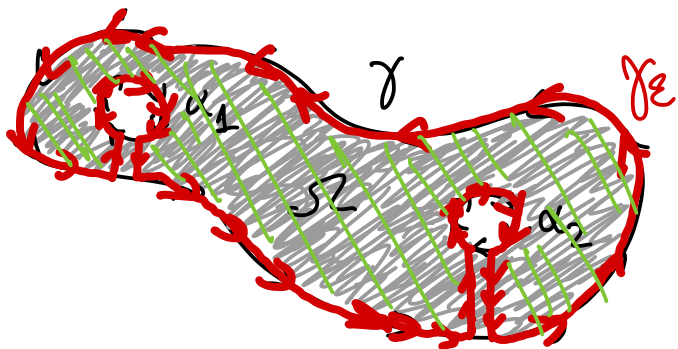
What if  $\Omega$  is not simply-connected?



$\gamma$ : outer boundary curve  
 (counterclockwise)

$\alpha_i$ : inner boundary curves  
 (clockwise)

By Green's Theorem applied to  $\gamma_{\epsilon}$  and  $\Omega_{\epsilon}$ :



$$\int_{\gamma_{\epsilon}} \vec{F} \cdot d\gamma_{\epsilon} = \iint_{\Omega_{\epsilon}} \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} dA$$

Note:  $\gamma_{\epsilon}$  bounds a simply connected region  $\Omega_{\epsilon}$

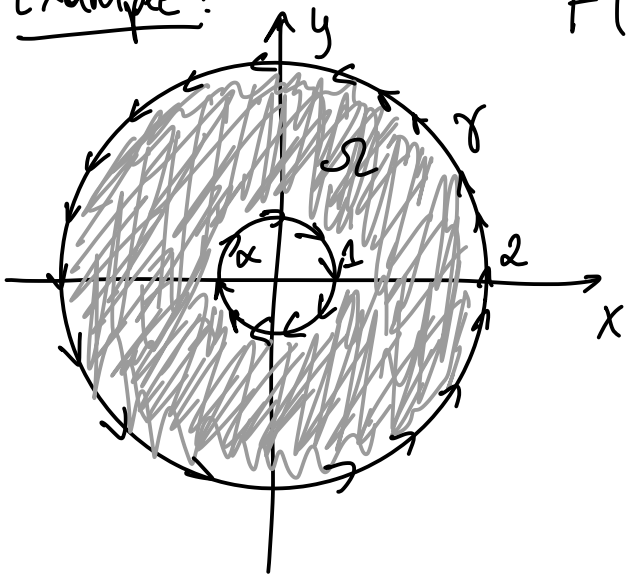
As  $\epsilon \downarrow 0$ , taking limits of the above, we get

$$\int_{\gamma} \vec{F} dy + \int_{\alpha_1} \vec{F} dx_1 + \int_{\alpha_2} \vec{F} dx_2 = \iint_{\Omega} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$$

or as many as the interior boundary curves  $\alpha_i$ , i.e.,

$$\sum_i \int_{\alpha_i} \vec{F} dx_i$$

Example:



$$\vec{F}(x,y) = (x+y, 1) = (M, N)$$

$$\int_{\gamma} \vec{F} dy + \int_{\alpha} \vec{F} dx = \iint_{\Omega} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$$

Compute RHS:

$$\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} = -1$$

$$\iint_{\Omega} \left( \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA = \iint_{\Omega} -1 dA = - \iint_{\Omega} 1 dA = - \text{Area}(\Omega)$$

$$= -(4\pi - \pi) = \boxed{-3\pi}$$

Compute LHS:

$$\gamma(t) = (2 \cos t, 2 \sin t), \quad t \in [0, 2\pi]$$

$$-\alpha(t) = (\cos t, \sin t), \quad t \in [0, 2\pi] \leftarrow \text{This is still counter-clockwise!}$$

$$\begin{aligned}
\int_{\gamma} \vec{F} \, d\gamma &= \int_0^{2\pi} \left\langle \underbrace{(2 \cos t + 2 \sin t, 1)}_{\vec{F}(\gamma(t))}, \underbrace{(-2 \sin t, 2 \cos t)}_{\gamma'(t)} \right\rangle dt \\
&= \int_0^{2\pi} -4 \cos t \sin t - 4 \sin^2 t + 2 \cos t \, dt \\
&= \int_0^{2\pi} \underbrace{-2 \sin 2t + 2 \cos t - 4 \sin^2 t}_{\substack{\text{this part integrates} \\ \text{to zero.}}} \, dt \\
&= -4 \int_0^{2\pi} \frac{1 - \cos 2t}{2} \, dt = -2(2\pi) = -4\pi
\end{aligned}$$

$$\begin{aligned}
\int_{-\alpha} \vec{F} \, d\alpha &= \int_0^{2\pi} \left\langle \underbrace{(\cos t + \sin t, 1)}_{\vec{F}(\alpha(t))}, \underbrace{(-\sin t, \cos t)}_{\alpha'(t)} \right\rangle dt \\
&= \int_0^{2\pi} -\cos t \sin t - \sin^2 t + \cos t \, dt \\
&= - \int_0^{2\pi} \frac{1 - \cos 2t}{2} \, dt = -\pi \quad \leftarrow \text{integrate to zero}
\end{aligned}$$

Thus  $\int_{\alpha} \vec{F} \, d\alpha = \pi$ .

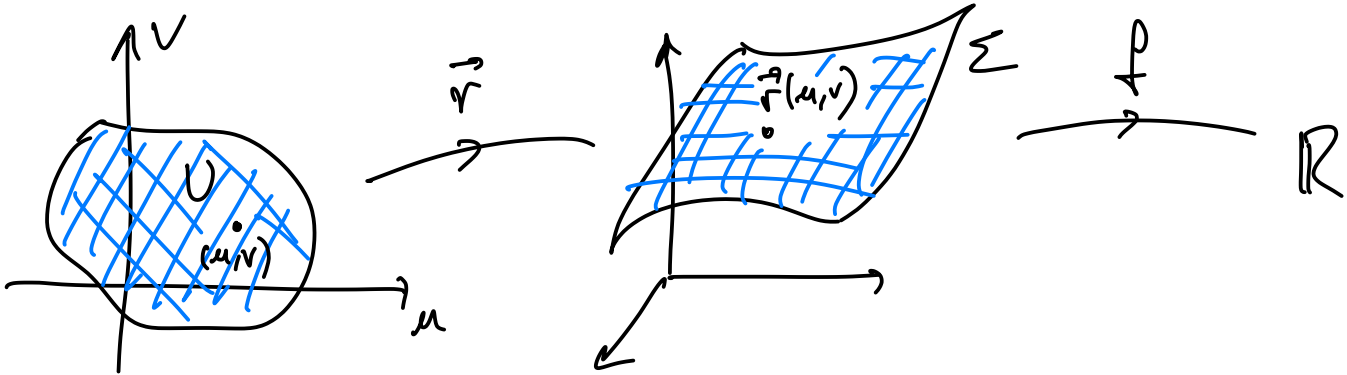
Altogether:  $\int_{\gamma} \vec{F} \, d\gamma + \int_{\alpha} \vec{F} \, d\alpha = -4\pi + \pi = \boxed{-3\pi}$

# A quick tour of more advanced Vector Calculus:

(Don't worry: there will be no questions on the final exam about this)

## Surface integrals:

$$\vec{r}: U \subset \mathbb{R}^2 \longrightarrow \Sigma \subset \mathbb{R}^3$$



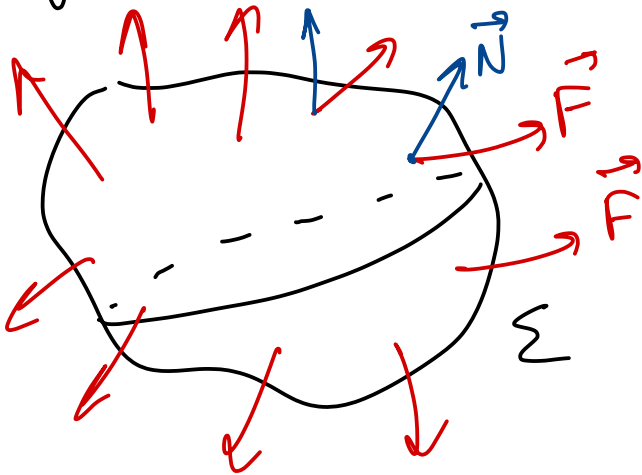
$$\vec{r}(u,v) = (x(u,v), y(u,v), z(u,v))$$

Surface integral

$$\iint_{\Sigma} f \, dS = \iint_U f(\vec{r}(u,v)) \cdot \left\| \frac{\partial \vec{r}}{\partial u} \times \frac{\partial \vec{r}}{\partial v} \right\| \, dA$$

Flux integral:

$$\iint_{\Sigma} \langle \vec{F}, \vec{N} \rangle \, dS = \text{flux of } \vec{F} \text{ across } \Sigma.$$



$\vec{F}$  vector field  
 $\vec{N}$  unit vector field

# Gauss Theorem / Divergence Theorem:

$$\iint_{\Sigma} \langle \vec{F}, \vec{N} \rangle dS = \iiint_{\Omega} \operatorname{div} \vec{F} dV$$

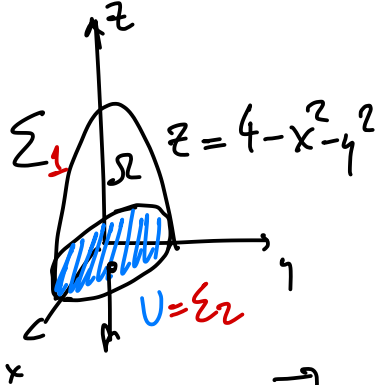
flux of  $\vec{F}$   
across  $\Sigma$

triple integral of  $\operatorname{div} \vec{F}$   
inside the domain  $\Omega$   
bounded by  $\Sigma$ .



$$\begin{aligned} \operatorname{div} \vec{F} &= \nabla \cdot \vec{F} = \left( \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \cdot (M, N, P) \\ &= \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} + \frac{\partial P}{\partial z} \end{aligned}$$

Example. Compute the flux of  $\vec{F}(x, y, z) = (2x, 2y, z)$  across the surface  $\Sigma$  in the picture.



$$\iint_{\Sigma} \langle \vec{F}, \vec{N} \rangle dS = \iiint_{\Omega} \operatorname{div} \vec{F} dV = \iiint_{\Omega} 5 dV = \dots$$

$\Sigma = \Sigma_1 \cup \Sigma_2$

$$\operatorname{div} \vec{F} = \frac{\partial}{\partial x} (2x) + \frac{\partial}{\partial y} (2y) + \frac{\partial}{\partial z} (z) = 2 + 2 + 1 = 5.$$

Parametrize  $\Omega$ :

$$0 \leq r \leq 2$$

$$0 \leq \theta \leq 2\pi$$

$$0 \leq z \leq 4 - r^2$$

$$\dots = \int_0^{2\pi} \int_0^2 \int_0^{4-r^2} 5 dz r dr d\theta =$$

$$= 10\pi \int_0^2 (zr) \Big|_0^{4-r^2} dr = 10\pi \int_0^2 (4-r^2)r dr$$

$$= 10\pi \int_0^2 4r - r^3 dr = 10\pi \left( 2r^2 - \frac{r^4}{4} \right) \Big|_0^2 = 10\pi(8-4) = \boxed{40\pi}$$

Note: We should also account for the "bottom" part of the boundary of  $\Omega$ , denoted  $\Sigma_2$  above:

$$\iint_{\Sigma_2} \langle \vec{F}, \vec{N} \rangle dS = \iint_U \langle (2r \cos \theta, 2r \sin \theta, 0), (0, 0, -1) \rangle dA = 0$$

$$\vec{F}(x, y, z) = (2x, 2y, z)$$

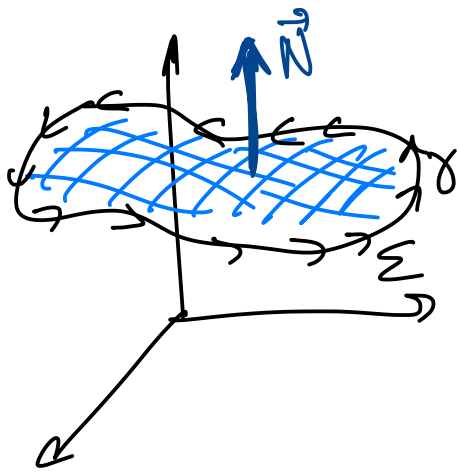
$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta \\ z = 0 \end{cases} \quad \begin{cases} r \in [0, 2] \\ \theta \in [0, 2\pi] \end{cases}$$

$$= \iint_U 0 dA = \boxed{0}$$

There's no flux through the bottom part  $\Sigma_2$ !

Stokes Theorem:

$\Sigma \subset \mathbb{R}^3$  surface with unit normal  $\vec{N}$   
and boundary  $\gamma$   
 $\vec{F}$  vector field



$$\int_{\gamma} \vec{F} d\gamma = \iint_{\Sigma} \langle \nabla \times \vec{F}, \vec{N} \rangle dS$$

flux of  $\text{curl } \vec{F} = \nabla \times \vec{F}$   
across  $\Sigma$

## General Version of Stokes Theorem:

$$\int_{\Omega} dw = \int_{\partial\Omega} w$$

$\Omega \subset \mathbb{R}^n$  open subset  
 $w$  is an  $(n-1)$ -form.